

## ON ANAND CONJECTURE

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### 1. Introduction

Unitons are harmonic maps  $\varphi : S^2 \rightarrow U(N)$ , which are important object in mathematics and physics. The construction of unitons was investigated from different perspectives by Uhlenbeck, Valli, Segal, Wood, Guest-Ohnita, Gu-Hu, Anand, and Burstall-Guest. For recent progress in this area see [12], [2], [11], [5] and their references.

In [1,2], Anand showed that unitons are equivalent to holomorphic ‘uniton bundles’, with energy corresponding to the bundles’ second Chern class. Using monad representation, he obtained a simple formula for the unitons and proved that 2-unitons have normalized energy at least four. This bound is sharp. In a personal communication with the author, Anand proposed the following conjecture:

**Conjecture (Finite Gap):**  $m$ -unitons have energy at least  $m^2$ .

The purpose of this note is to give a positive answer to this conjecture. It is known that harmonic maps of simply connected Riemann surfaces into  $U(N)$  can be lifted into holomorphic maps into the based loop group  $\Omega U(N)$ , which are called extended harmonic maps. Unitons have algebraic extended harmonic maps  $\Phi : S^2 \rightarrow \Omega_{\text{alg}} U(N)$  ( see [20] for details). Our idea is to deform a given extended harmonic map into a simpler one, which allows us to confirm the above conjecture of Anand.

### 2. Harmonic maps and extended harmonic maps

Let  $\varphi : R^2 \rightarrow U(N)$  be a smooth map. Define  $A = \frac{1}{2}\varphi^{-1}d\varphi$ , and decompose it as follows:

$$A = A_z dz + A_{\bar{z}} d\bar{z} \quad \text{with } (A_z)^* = -A_{\bar{z}},$$

where  $z = x + iy \in R^2$ . Then  $A$  satisfies the Maurer-Cartan equation:

$$(1) \quad \bar{\partial}A_z - \partial A_{\bar{z}} + 2[A_{\bar{z}}, A_z] = 0.$$

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In the case of maps into a matrix group, with the standard (left-invariant) metric, the energy takes the form

$$\begin{aligned} E(\varphi) &= \frac{1}{2} \int_{R^2} (|\varphi^{-1} \frac{\partial \varphi}{\partial x}|^2 + |\varphi^{-1} \frac{\partial \varphi}{\partial y}|^2) dx \wedge dy \\ (2) \qquad &= -4i \int \text{tr} A_z A_{\bar{z}} dz d\bar{z}. \end{aligned}$$

It is known that  $\varphi$  is harmonic if and only if  $A$  is co-closed; i.e.,

$$(3) \qquad d^* A = \frac{\partial}{\partial \bar{z}} A_z + \frac{\partial}{\partial z} A_{\bar{z}} = 0.$$

Unions are harmonic maps from  $S^2$  to  $U(N)$ . From [19], we know that harmonic maps from  $R^2 \rightarrow U(N)$  extend to  $S^2$  iff they have finite energy, and that such maps are always smooth. In the following, we will use this fact and work in terms of coordinates  $x$  and  $y$  on  $R^2$ .

The Lax pair of the harmonic maps is ([20]):

$$(4) \qquad \begin{cases} \bar{\partial} \Phi = (1 - \lambda) \Phi A_{\bar{z}}, \\ \partial \Phi = (1 - \lambda^{-1}) \Phi A_z, \end{cases}$$

whose integrability conditions are (1) and (3). A nondegenerated  $N \times N$  matrix solution  $\Phi(\lambda) := \Phi_\lambda$  to (4) is called extended solution or extended harmonic map. It can be normalized so that

$$(5) \qquad \Phi_1 = I, \quad \Phi_\lambda \Phi_{\sigma(\lambda)}^* = I.$$

where  $\sigma(\lambda) = (\bar{\lambda})^{-1}$ . Moreover, we have  $\Phi_{-1} = Q\varphi$  for some constant  $Q \in U(N)$ .

Let  $\Omega U(N) = \{\gamma : S^1 \rightarrow U(N) | \gamma \text{ is smooth, } \gamma(1) = I\}$  be the group of smooth based loops in  $U(N)$ . By (5), we can regard an extended harmonic map as a smooth map into  $\Omega U(N)$ . Denote by  $\Omega_{\text{alg}} U(N)$  the subgroup of  $\Omega U(N)$  consisting of algebraic loops. One of the main results of Uhlenbeck [20] concerning harmonic maps  $\varphi : S^2 \rightarrow U(N)$  is :

**Theorem 2.1[20].** *Let  $\varphi : S^2 \rightarrow U(N)$  be harmonic. Then there exists an algebraic extended harmonic map  $\Phi : M \rightarrow \Omega_{\text{alg}} U_n$  of  $\varphi$ .*

It is easy to see that the algebraic extended harmonic map  $\Phi$  can be expressed as the following type after multiplication on the left by a suitable constant loop  $\gamma \in \Omega U(N)$ :

$$(6) \qquad \Phi = \sum_{\alpha=0}^m \lambda^\alpha T_\alpha.$$

According to [20], we call the harmonic map  $\varphi : S^2 \rightarrow U(N)$  and the extended harmonic map  $\Phi : S^2 \rightarrow \Omega U(N)$  a uniton and an extended uniton respectively. We say that  $\varphi$  or  $\Phi$  has finite uniton number  $m$ . The minimal such number  $m$  is called the minimal uniton number of  $\varphi$  or  $\Phi$ . It was shown in [20] that an extended uniton  $\Phi$  can be normalized further so that

$$(7) \quad V_0(\Phi) := \text{Span}\{Im_{z \in S^2} T_0(z)\} = C^N,$$

and if an extended uniton  $\Phi$  satisfies the condition (7) then  $m$  is just the minimal uniton number. An extended uniton  $\Phi$  with the property (7) will be called the standard extended uniton of  $\varphi$ .

**lemma 2.2.** *Let  $\Phi_\lambda = \sum_{\alpha=0}^m \lambda^\alpha T_\alpha : S^2 \rightarrow \Omega U(N)$  be an extended uniton. Then  $ImT_0 : S^2 \rightarrow G_k(C^N)$  is a holomorphic map, where  $\text{rank}(ImT_0) = k$ . Furthermore,  $\Phi_\lambda$  is standard if and only if  $ImT_0$  is a linearly full holomorphic curve in  $G_k(C^N)$ . Similarly,  $ImT_m$  is an anti-holomorphic map.*

*Proof.* We have two unitary connections on the trivial bundle  $\underline{C}^N : d_A = d + A$  and the trivial connection  $d$ . They determine their corresponding Koszul-Malgrange holomorphic structures on  $\underline{C}^N$  by [15]. From (4), we have

$$\bar{\partial}T_0 = T_0 A_{\bar{z}}$$

which means that  $T_0$  is a global holomorphic section of  $\text{Hom}[(\underline{C}^N, d_A), (\underline{C}^N, d)]$ . Thus  $ImT_0$  is a holomorphic subbundle of  $(\underline{C}^N, d)$ , i.e.,  $ImT_0$  is a holomorphic map to  $G_k(C^N)$ .  $\Phi_\lambda$  is standard if and only if it satisfies (7), which means that  $ImT_0$  is a linearly full holomorphic map. By (4), we have also

$$\partial T_m = T_m A_z$$

We can prove similarly that  $ImT_m$  is an anti-holomorphic map.  $\square$

We introduce various fundamental loop groups as follows:

$$\begin{aligned} \Lambda GL(N, C) &= \{\gamma : S^1 \rightarrow GL(N, C) \mid \gamma \text{ is smooth}\}, \\ \Lambda^+ GL(N, C) &= \{\gamma \in \Lambda GL(N, C) \mid \gamma \text{ extends continuously to} \\ &\quad \text{a holomorphic map } D^+ \rightarrow GL(N, C)\} \end{aligned}$$

where  $D^+ = \{\lambda \mid |\lambda| < 1\}$ . We recall the Iwasawa decomposition:

**Theorem 2.3**[18]. *Each  $\gamma$  in  $\Lambda GL(N, C)$  can be uniquely decomposed as*

$$\gamma = \gamma_u \gamma_+,$$

where  $\gamma_u \in \Omega U(N)$  and  $\gamma_+ \in \Lambda^+ GL(N, C)$ .

By the theorem, we have an identification

$$(8) \quad \Omega U(N) \cong \Lambda GL(N, C) / \Lambda^+ GL(N, C)$$

and thus, a natural action of the complex group  $\Lambda GL(N, C)$  on the complex manifold  $\Omega U(N)$  ([13]): if  $\gamma \in \Lambda GL(N, C)$  and  $\delta \in \Omega U(N)$ , then

$$(9) \quad \gamma \cdot \delta = (\gamma\delta)_u.$$

Let us choose a maximal torus  $T$  of  $U(N)$  and a fundamental Weyl chamber  $W_0$  in the Lie algebra  $t$  of  $T$ . Denote by  $I'$  the intersection of the integer lattice  $I = (2\pi)^{-1} \exp^{-1}(e) \cap t$  with  $W_0$ . Then we know that  $I' (\ni \xi)$  parametrizes all the conjugacy classes of homomorphisms  $\gamma_\xi : S^1 \rightarrow U(N)$ . By [18], we have the following Bruhat decomposition :

$$(10) \quad \Omega_{\text{alg}} U(N) = \sqcup_{\xi \in I'} \Lambda_{\text{alg}}^+ GL(N, C) \cdot \gamma_\xi,$$

where  $\Lambda_{\text{alg}}^+ GL(N, C) = \Lambda_{\text{alg}} GL(N, C) \cap \Lambda^+ GL(N, C)$ .

There is a beautiful Morse theoretic explanation of these decompositions. Let  $H : \Omega U(N) \rightarrow R$  be the usual energy functional on paths,  $H(\gamma) = \int_{S^1} |\gamma'|^2$ . This is a Morse-Bott function. The critical manifolds are the conjugacy classes of homomorphisms. We write

$$(11) \quad \Omega_\xi = \{g\gamma_\xi g^{-1} \mid g \in U(N)\}$$

for the conjugacy class of the homomorphism  $\gamma_\xi$ . Denote by  $U_\xi$  the unstable manifold of  $\Omega_\xi$  with respect to the flow of  $-gradH$  (i.e.,  $U_\xi$  is the stable manifold of  $\Omega_\xi$  with respect to the flow of  $gradH$ ). Then we have ([17])

$$(12) \quad U_\xi = \Lambda_{\text{alg}}^+ GL(N, C) \cdot \gamma_\xi.$$

In [17], the energy functional of loops  $H$  is viewed as a Hamiltonian function associated to a symplectic action of the group  $S^1$  on  $\Omega U(N)$ , which is given by

$$(13) \quad \mu \cdot \gamma(\lambda) = \gamma(\mu\lambda)\gamma(u)^{-1}, \quad \mu \in S^1, \gamma \in \Omega U(N).$$

It was observed by Terng (see sect.7 of [20]) that this action preserves the extended harmonic equation (4). We have the following action of  $C^*$  on  $\Omega_{\text{alg}} U(N)$ :

$$(14) \quad \mu \cdot \gamma(\lambda) = \gamma(\mu\lambda)\Lambda^+ GL(N, C), \quad \mu \in C^*, \gamma \in \Omega_{\text{alg}} U(N),$$

which is the complexification of the  $S^1$ -action. It turns out that the flow line of  $gradH$  starting at a point  $\gamma \in \Omega_{\text{alg}} U(N)$  is given by the action of subsemigroup  $(0, 1]$  of  $C^*$ .

For an extended unition  $\Phi : S^2 \rightarrow \Omega_{\text{alg}} U(N)$ , we cannot conclude that  $\Phi(S^2) \subseteq U_\xi$  for some  $\xi \in I'$  in general. However, we have

**proposition 2.4[5].** *Let  $\Phi : S^2 \rightarrow \Omega_{\text{alg}} U(N)$  be an extended unition. Then there exists some  $\xi \in I'$ , and some discrete subset  $D$  of  $S^2$ , such that  $\Phi(S^2 - D) \subseteq U_\xi$ .*

### 3. Deformations of $S^1$ -invariant unitons

In the classification theory of unitons in [5], a fundamental role is played by the following special extended unitons:

**Definition 3.1.** An  $S^1$ -invariant extended uniton is an extended uniton  $\Phi$  such that  $Im\Phi \subseteq \Omega_\xi$ , for some  $\xi \in I'$ ; or equivalently  $\Phi$  is invariant under the  $S^1$ -action (13). In this case,  $\Phi_{-1}$  is also called  $S^1$ -invariant uniton.

By [20], we know that  $\Phi$  is an  $S^1$ -invariant extended uniton into  $\Omega U(N)$  if and only if it can be expressed as

$$(15) \quad \Phi = \sum_{\alpha=0}^k \Pi_\alpha \lambda^\alpha$$

where  $f_\Phi := (\Pi_0, \Pi_1, \dots, \Pi_k)$  is a super-horizontal holomorphic map into a complex flag manifold  $F = F(n_0, n_1, \dots, n_k)$  with  $n_\alpha = rank \Pi_\alpha$  ( $\alpha = 0, 1, \dots, k$ ) and  $N = n_0 + n_1 + \dots + n_k$ . We say that  $f_\Phi$  is standard if the extended uniton  $\Phi$  is standard.

In the following, we will investigate the deformation of the  $S^1$ -invariant harmonic map. We know that  $F(n_0, \dots, n_k)$  can be identified with the complex homogeneous space  $U(N)/U(n_0) \times \dots \times U(n_k)$ . A point  $B$  of  $F$  may be viewed as a  $(k+1)$ -tuple of orthogonal subspaces  $B_0 \oplus \dots \oplus B_k$  where  $\dim(B_i) = n_i$ , or, equivalently, as a flag

$$0 \subset V_0 = B_0 \subset V_1 = B_0 + B_1 \subset \dots \subset V_k = B_0 + B_1 + \dots + B_k = C^N.$$

Let  $\{V_0, \dots, V_k\}$  denote the flag corresponding  $(B_0, \dots, B_k)$ . An element of  $F(n_0, \dots, n_k)$  is written as  $B = (B_0, B_1, \dots, B_k)$  or  $\{V_0, V_1, \dots, V_k\}$ . Set  $J = \{0, 2, \dots, 2\lfloor \frac{k}{2} \rfloor\}$  and  $J^c = \{0, 1, \dots, k\} \setminus J$ . Define the projection

$$(16) \quad \pi : F(n_0, n_1, \dots, n_k) \ni (B_0, B_1, \dots, B_k) \rightarrow \bigoplus_{j \in J^c} B_j \in Gr_t(C^N)$$

where  $t = \sum_{j \in J^c} n_j$ . It is known that  $\pi$  is the so called canonical twistor fibration (see [6] or [9]).

Define the following tautological complex vector bundles for  $B \in F(n_0, n_1, \dots, n_k)$  :  $(V_0)_B = B_0, (V_1)_B = B_0 \oplus B_1, \dots, (V_k)_B = \bigoplus_{i=0}^k B_i$ . A map  $f : S^2 \rightarrow F(n_0, n_1, \dots, n_k)$  is called super-horizontal holomorphic map if

$$(17) \quad \begin{cases} \partial' C^\infty(f^{-1}V_i) \subset C^\infty(f^{-1}V_i), \\ \partial' C^\infty(f^{-1}V_i) \subset C^\infty(f^{-1}V_{i+1}), \quad i = 0, 1, \dots, k, \end{cases}$$

where we set  $V_{k+1} = C^N$ .

Hence a uniton  $\varphi : S^2 \rightarrow Gr_t(C^N)$  is  $S^1$ -invariant if and only if there exists a super-horizontal holomorphic map  $f_\Phi : S^2 \rightarrow F$  such that  $\varphi = \pi \circ f_\Phi$ . The group

$GL(N, C)$  acts on  $F(n_0, n_1, \dots, n_k)$  in the natural way : for  $A \in GL(N, C)$  and  $V_0 \subset V_1 \subset \dots \subset V_k$ , we have  $AV_0 \subset AV_1 \subset \dots \subset AV_k$ . It is easy to see that this action preserves the super- horizontality and holomorphicity conditions. Thus it induces an action on the  $S^1$ -invariant untions.

Let  $G = U(N)$  and  $g = u(N)$  denote the unitary group and its Lie algebra. Then we can regard  $F = F(n_0, n_1, \dots, n_k)$  as an “Ad-orbit” of  $U(N)$  as follows: If we let  $(\widehat{B}_0, \widehat{B}_1, \dots, \widehat{B}_k)$  be a fixed element of  $F(n_0, n_1, \dots, n_k)$  and set  $\xi = \sqrt{-1} \sum_{i=0}^k \pi_{W_i}$  where  $W_i = \bigoplus_{\alpha=0}^i \widehat{B}_\alpha$  for  $i = 0, 1, \dots, k$ , then we have  $F(n_0, n_1, \dots, n_k) \cong Ad(G)\xi$ . Here  $\pi_V$  denotes the Hermitian projection in  $C^N$  onto a complex subspace  $V$ .

Fix an element  $L \in CP^{N-1}$  and put  $Q = \sqrt{-1}\pi_L \in u(N)$ . For a  $X = \sqrt{-1} \sum_{i=0}^k \pi_{V_i} \in Ad(G)\xi$ , where  $(V_0, V_0^\perp \cap V_1, \dots, V_{k-1} \cap V_k^\perp) \in F(n_0, n_1, \dots, n_k)$ , we define the height function  $h^Q : Ad(G)\xi \rightarrow R$  by

$$(18) \quad h^Q(X) = (X, Q).$$

Here  $(,)$  is an  $Ad(G)$ -invariant inner product on  $g$ . Then it is known that  $h^Q$  is a Morse-Bott function. Let  $gradh^Q$  be the gradient of  $h^Q$  with respect of the Kähler metric. The following fact is due to Frankel [16] : the flow of  $-gradh^Q$  is given by the natural action of  $\{\exp \sqrt{-1}tQ\}$ . A point  $(V_0, V_0^\perp \cap V_1, \dots, V_{k-1}^\perp \cap V_k)$  is a critical point if and only if  $[\sqrt{-1} \sum_{i=0}^k \pi_{V_i}, \sqrt{-1}\pi_L] = 0$ , i.e., the line  $L$  is contained in one of  $V_0, V_0^\perp \cap V_1, \dots, V_{k-1}^\perp \cap V_k$ .

**lemma 3.1.** *There are  $k + 1$  connected no-degenerate critical manifolds of  $h^Q$  :*

$$F_0 = \{\{V_0, V_1, \dots, V_k\} \in F(n_0, n_1, \dots, n_k) | L \subset V_0\} \cong F(n_0 - 1, n_1, \dots, n_k),$$

$$F_1 = \{\{V_0, V_1, \dots, V_k\} \in F(n_0, n_1, \dots, n_k) | L \subset V_0^\perp \cap V_1\} \cong F(n_0, n_1 - 1, \dots, n_k),$$

.....

$$F_k = \{\{V_0, V_1, \dots, V_k\} \in F(n_0, n_1, \dots, n_k) | L \subset V_{k-1}^\perp \cap V_k\} \cong F(n_0, n_1, \dots, n_k - 1).$$

We set  $G_Q = \{A \in GL(N, C) | A(L) = L\}$ . It is known that the stable manifold for a non-degenerate critical manifold  $N$  is given by  $S^Q(N) = G_Q X$  for  $X \in N$ . In this case, we have

**lemma 3.2.** *The corresponding stable manifolds are*

$$S^Q(F_0) = F_0,$$

$$S^Q(F_1) = \{\{V_0, V_1, \dots, V_k\} \in F(n_0, n_1, \dots, n_k) | V_0 \cap L = \{0\}, L \subset V_1\},$$

.....

$$S^Q(F_k) = \{\{V_0, V_1, \dots, V_k\} \in F(n_0, n_1, \dots, n_k) | V_{k-1} \cap L = \{0\}, L \subset V_k\}.$$

Let  $\varphi : S^2 \rightarrow Gr_t(C^N)$  be a  $S^1$ -invariant harmonic map, and  $f_\Phi : S^2 \rightarrow F$  be a super-horizontal holomorphic map corresponding to  $\varphi$ . If  $f_\Phi(S^2) \subset S^Q(F_k)$ ,

then  $\{(\exp \sqrt{-1}tQ) \circ f_\Phi\}$  provides a continuous deformation to a super-horizontal holomorphic map into  $F_k(\cong F(n_0, n_1, \dots, n_k - 1))$ . We shall show that there exists some  $L \in CP^{N-1}$  such that  $f_\Phi(S^2) \subset S^Q(F_k)$  if  $n_k \geq 2$ .

Set  $Y = \{L \in CP^{N-1} | f_\Phi(z) \notin S^Q(F_k) \text{ for some } z \in S^2\}$ . Then we have  $Y = \{L \in CP^{N-1} | L \subset V_{k-1}(z) \text{ for some } z \in S^2\}$ . It suffices to show that  $Y$  cannot be equal to  $CP^{N-1}$ .

Fix  $z \in S^2$ , we set  $Y_z = \{L \in CP^{N-1} | L \subset V_{k-1}(z)\}$ . It is easy to see that

$$(19) \quad \dim_C Y_z \leq \sum_{i=0}^{k-1} n_i - 1.$$

Since  $\dim_C S^2 = 1$ , we have

$$(20) \quad \dim_C Y \leq \sum_{i=0}^{k-1} n_i.$$

It follows that, if  $n_k \geq 2$ , the space  $Y$  cannot be equal to  $CP^{N-1}$ . It suffices to choose  $L \in CP^{N-1} \setminus Y$ . Hence, by induction, we obtain a map whose image lies in  $F(n_0, n_1, \dots, n_{k-1}, 1)$ . By repeating the argument with  $f^* = \{V_k^\perp, \dots, V_0^\perp, C^n\} \in F(1, n_{k-1}, \dots, n_0)$ , we can similarly deform  $f_\Phi$  into  $\{\{V_0, V_1, \dots, V_k\} \in F(n_0, n_1, \dots, n_{k-1}, 1) | C^{n_0-1} \subseteq V_0\} \cong F(1, n_1, \dots, n_{k-1}, 1)$ . Thus we obtain the following proposition.

**proposition 3.3.** *If  $n \geq 3$ , then any super-horizontal holomorphic map  $f_\Phi$  into  $F(n_0, n_1, \dots, n_k)$  can be deformed continuously through super-horizontal holomorphic maps to a super-horizontal holomorphic map  $f_{\tilde{\Phi}}$  into  $F(1, n_1, \dots, n_{k-1}, 1)$ . Furthermore, if  $f_\Phi$  is standard, then  $f_{\tilde{\Phi}}$  is also standard in its corresponding space.*

*Proof.* We only need to prove the second part of the proposition. Assume that  $f_\Phi$  is standard, i.e., the corresponding extended uniton  $\Phi = \sum_{\alpha=0}^k \Pi_\alpha \lambda^\alpha$  satisfies the following condition

$$V_0(\Phi) = \text{Span}\{Im_{z \in S^2} \Pi_0(z)\} = C^N$$

Let  $\Phi^t = \sum_{\alpha=0}^k \Pi_\alpha^t \lambda^\alpha$  be the continuous deformation obtained in the above. It is clear that  $\dim V_0(\Phi^t) = \text{const}$ . So,  $V_0(\Phi^t) = C^N$ . In particular, we have

$$V_0(\Phi^\infty) = C^N.$$

The above deformation shows that  $f_{\Phi^\infty}(S^2) \subset \{\{V_0, V_1, \dots, V_k\} \in F(n_0, n_1, \dots, n_k) | C^{n_0-1} \subseteq V_0, C^{n_k-1} \subseteq V_{k-1}^\perp \cap V_k\} \stackrel{i}{\cong} F(1, n_1, \dots, n_{k-1}, 1)$ , where  $i$  denotes the natural isomorphism. Let  $f_{\tilde{\Phi}} = i \circ f_{\Phi^\infty} : S^2 \rightarrow F(1, n_1, \dots, n_{k-1}, 1)$ . We see that

$$V_0(\tilde{\Phi}) = C^{\tilde{N}}$$

where  $\tilde{N} = N - (n_0 + n_k - 2)$ . Hence  $\tilde{\Phi}$  is a standard extended  $U(\tilde{N})$ -uniton, i.e.,  $f_{\tilde{\Phi}}$  is standard.  $\square$

We know from [8] that any harmonic map  $S^2 \rightarrow CP^{N-1}$  corresponds to a super-horizontal holomorphic map into  $F(n_0, 1, n_2)$  with  $n_0 + 1 + n_2 = N$ . Hence we obtain:

**Corollary 3.4[13].** *If  $N \geq 3$ , then any harmonic map  $\varphi : S^2 \rightarrow CP^{N-1}$  can be deformed continuously through harmonic maps to a harmonic map  $\psi : S^2 \rightarrow CP^2$ .*

#### 4. Anand's conjecture

We now give an application of the above deformations for unitons.

**Theorem 4.1.** *Let  $\varphi : S^2 \rightarrow U(N)$  be a harmonic map with minimal uniton number  $m$ . Then*

$$(21) \quad E(\varphi) \geq 4\pi m^2.$$

Furthermore, the equality holds if and only if  $\varphi$  can be deformed continuously to a harmonic map  $\psi$  which is one of the following two cases :

- (i)  $\psi : S^2 \rightarrow CP^1(\subset U(2) \subset U(N))$  is a totally unramified holomorphic map ;
- (ii)  $\psi : S^2 \rightarrow Gr_2(C^3)$  is a harmonic map with minimal uniton number  $m = 2$  and  $\psi^\perp : S^2 \rightarrow CP^2(\subset U(3) \subset U(N))$  is a totally unramified, not  $\pm$ holomorphic , full harmonic map.

*Proof.* Let  $\Phi = \sum_{\alpha=0}^m T_\alpha \lambda^\alpha$  be the standard extended harmonic map of  $\varphi$ . By Proposition 2.3 we can express  $\Phi$  on  $M - D$  as follows:

$$(22) \quad \Phi = A\gamma_\xi B$$

where  $A, B : M - D \rightarrow \Lambda_{\text{alg}}^+ GL_n(C)$  and  $\gamma_\xi$  is the geodesic given by

$$\gamma_\xi(\lambda) = \text{diag}(\lambda^m, \dots, \lambda^m, \dots, \lambda, \dots, \lambda, 1, \dots, 1).$$

It is easy to see that

$$(23) \quad T_0 = A(0)\gamma_\xi(0, \dots, 1, \dots, 1)B(0).$$

on  $M - D$ . Now we deform the extended harmonic map  $\Phi$  by the gradient flow of the Morse-Bott function  $H : \Omega G \rightarrow R$ . From the description of this flow in Sec.2, we have  $\hat{\Phi} = \lim_{t \rightarrow +\infty} \Phi^t$ , where  $\Phi^t$  is the extended harmonic map  $e^{-t} \cdot \Phi$ . The points in  $D$  are the possible bubble points in the deformation. If bubbling occurs, then  $E(\varphi) > E(\hat{\Phi}_{-1})$ . So, in general , we have  $E(\varphi) \geq E(\hat{\Phi}_{-1})$ . Since  $Im\Psi \subseteq \Omega_\xi$ , the extended uniton  $\Psi$  is an  $S^1$ -invariant uniton. By [20]  $\hat{\Phi}$  has the following expression:

$$(24) \quad \hat{\Phi} = \sum_{\alpha=0}^m \Pi_\alpha \lambda^\alpha$$

where  $(\Pi_0, \Pi_1, \dots, \Pi_m)$  is a super-horizontal holomorphic map into  $F(n_0, n_1, \dots, n_m)$  with  $n_\alpha = \text{rank}(\Pi_\alpha)$  ( $\alpha = 0, 1, \dots, m$ ). From the definition of the action  $e^{-t} \in C^*$  on extended uniton and (22), we have

$$\begin{aligned} \widehat{\Phi} &= A(0) \cdot \gamma_\xi \\ &= A(0)\gamma_\xi C_0. \end{aligned}$$

We see that  $\pi_0 = A(0)\text{diag}(0, \dots, 0, 1, \dots, 1)C_0$ . It follows that

$$\begin{aligned} \text{Im}\pi_0 &= \text{Im}A(0)\text{diag}(0, \dots, 0, 1, \dots, 1) \\ &= \text{Im}T_0. \end{aligned}$$

Thus  $\widehat{\Phi}$  is a standard extended uniton with minimal uniton number  $m$ .

By Proposition 3.3,  $\widehat{\Phi}$  can be deformed continuously to an  $S^1$ -invariant extended harmonic map  $\Psi$  which corresponding a standard super-horizontal holomorphic map into  $F(1, n_1, \dots, n_{m-1}, 1)$ . Hence  $E(\widehat{\Phi}_{-1}) = E(\Psi_{-1})$  and the  $S^1$ -invariant extended harmonic map  $\Psi$  is given by

$$(25) \quad \Psi = \sum_{\alpha=0}^m (\pi_\alpha - \pi_{\alpha-1})\lambda^\alpha$$

with  $\text{rank}(\pi_0) = 1$ . From Lemma 2.2, we see that  $\psi_0 := \text{Im}(\pi_0) : S^2 \rightarrow CP^n$  is a full holomorphic map, where  $n = 1 + \sum_{i=1}^{m-1} n_i$ . The holomorphic map  $\psi_0$  generates the following harmonic sequence (see [4], [22]):

$$(26) \quad \psi_0 \xrightarrow{\partial_0} \psi_1 \xrightarrow{\partial_1} \dots \xrightarrow{\partial_{N-2}} \psi_n,$$

with  $\bigoplus_{j=0}^n \underline{\psi}_j = \underline{C}^{n+1}$ , where  $\partial$  denotes the fundamental collineation defined by Chern and Wolfson(see [7]). It is easy to see from the section10 of [20] that

$$(27) \quad \underline{\pi}_0 = \underline{\psi}_0, \underline{\pi}_1 = \bigoplus_{j=0}^{i_1} \underline{\psi}_j, \dots, \underline{\pi}_{m-1} = \bigoplus_{j=0}^{i_{m-1}} \underline{\psi}_j, \underline{\pi}_m = \underline{C}^{n+1},$$

with  $0 = i_0 < i_1 < \dots < i_{m-1} < i_m = n$ . It follows from Uhlenbeck's factorization and Valli's formula ([21]) that

$$(28) \quad \begin{aligned} E(\Psi_{-1}) &= 4\pi \sum_{\alpha=0}^m \text{deg } \pi_\alpha \\ &= 4\pi \sum_{\alpha=0}^{m-1} \text{deg} \left( \bigoplus_{j=0}^{i_\alpha} \underline{\psi}_j \right). \end{aligned}$$

Since  $\pi_\alpha$  corresponds to the  $i_\alpha$ -th osculating curve of the holomorphic curve  $\psi_0$ , we have by [4] that

$$\begin{aligned}
 \deg\left(\bigoplus_{j=0}^{i_\alpha} \psi_j\right) &= (i_\alpha + 1)(n - i_\alpha) + \frac{n - i_\alpha}{n + 1} \sum_{k=0}^{i_\alpha-1} (k + 1)r(\partial_k) \\
 &\quad + \frac{i_\alpha + 1}{n + 1} \sum_{k=i_\alpha}^{n-1} (n - k)r(\partial_k) \\
 (29) \qquad \qquad \qquad &\geq (i_\alpha + 1)(n - i_\alpha),
 \end{aligned}$$

where  $r(\partial_k)$  denotes the ramification index of  $\partial_k$ . It follows from (28) and (29) that  $E(\Psi_{-1}) \geq 4\pi mn$ . The minimal uniton number  $m \leq n$  by [20], so we obtain  $E(\varphi) \geq 4\pi m^2$ .

Set  $\psi = \Psi_{-1}$ . If the equality holds, then we see from (28) and (29) that

$$r(\partial_k) = 0$$

for  $k = 0, 1, \dots, n - 1$  and

$$(30) \qquad \qquad \qquad (i_\alpha + 1)(n - i_\alpha) = n$$

for  $\alpha = 0, 1, \dots, m - 1$ . It follows from (30) that either  $m = 1$  and  $n = 1$ , i.e.,  $\psi$  belongs to the case (i) or  $m = 2$ ,  $n = 2$  and

$$\begin{aligned}
 \psi &= \pi_0 - \pi_1 + \pi_2 \\
 &= \pi_0 + \pi_2 - \pi_1,
 \end{aligned}$$

i.e.,  $\psi$  belongs to the case (ii). Actually, we see via the Cartan embedding that  $\psi = \psi_0$  in the case (i); and  $\psi = \psi_0 \oplus \psi_2$ ,  $\psi^\perp = \psi_1$  in the case (ii). Conversely, if  $\psi$  is one of the cases (i) and (ii), it is easy to verify that the equality in (21) holds.  $\square$

*Remark 4.1.* The above proof shows a little more, namely that if  $\varphi : S^2 \rightarrow U(N)$  has a standard extended harmonic map  $\Phi = \sum_{\alpha=0}^m T_\alpha \lambda^\alpha$  with  $\max_{z \in S^2} \text{rank}(T_0) = 1$ , then  $E(\varphi) \geq 4\pi m(N - 1)$ . We can normalize the energy to be  $\tilde{E}(\varphi) = \frac{1}{4\pi} E(\varphi)$ . So, the Anand conjecture states that  $m$ -unitons have energy at least  $m^2$ .

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